

A Solar Stationary Type IV Radio Burst and its Radiation Mechanism

Hongyu Liu

Sep. 21, 2017 Jeju

Hongyu Liu
henryleo0@mail.ru

¹ Korea Astronomy and Space Science Institute, 776 Daedeokdae-ro, Yuseong-gu, Daejeon 34055, Korea



² University of Science and Technology, 217 Gajeong-ro, Yuseong-gu, Daejeon 34113, Korea



³ Shandong Provincial Key Laboratory of Optical Astronomy and Solar-Terrestrial Environment, and Institute of Space Sciences, Shandong University, Weihai, Shandong 264209, China



Outline

1 Introduction

2 Observations of Type IV Event
on 20110924

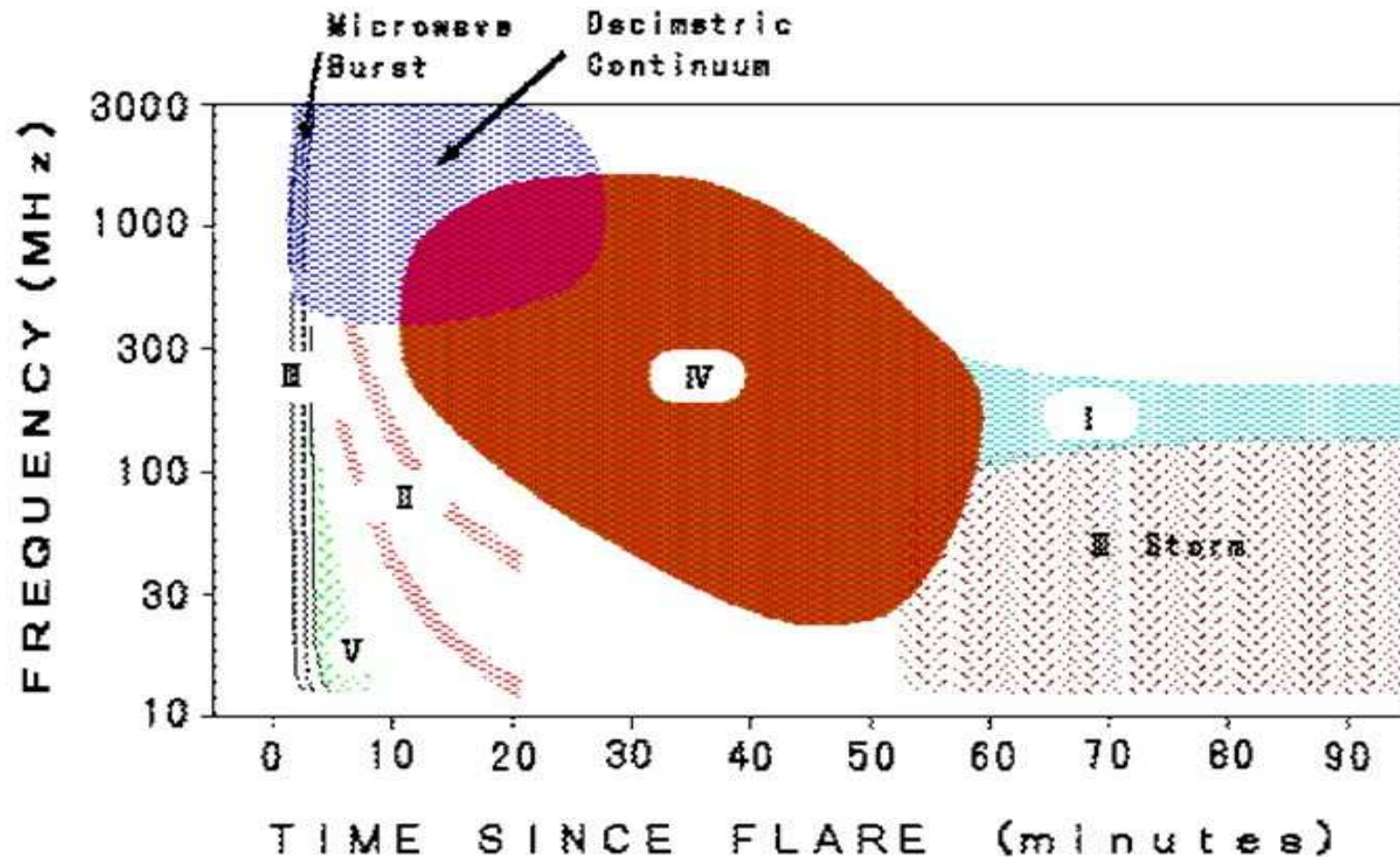
3 Summary&Radiation
Mechanism Interpretation

1 Introduction

- 1) Solar radio bursts and Type IV bursts
- 2) Important Physical parameters
- 3) Stationary Type IV general features
- 4) Gyrosynchrotron and plasma radiation
- 5) Solar radio spectrometer and radioheliograph
- 6) Previous similar study

What is Solar Radio Burst?

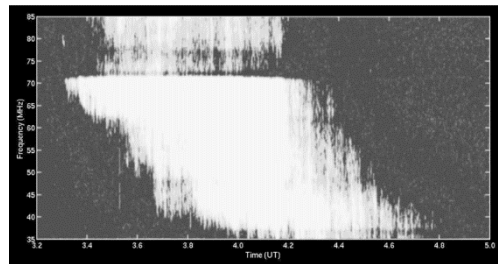
intense increase of radio frequency radiation induced by solar eruptions



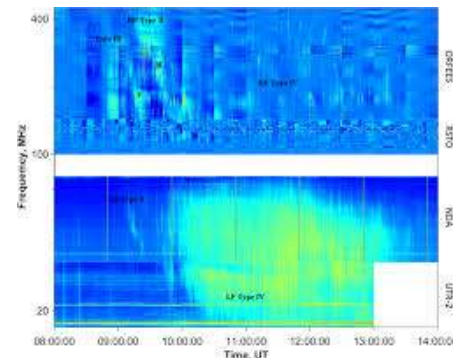
Type IV Radio Bursts

Broadband meter wavelength continuum which lasts for relatively long time.

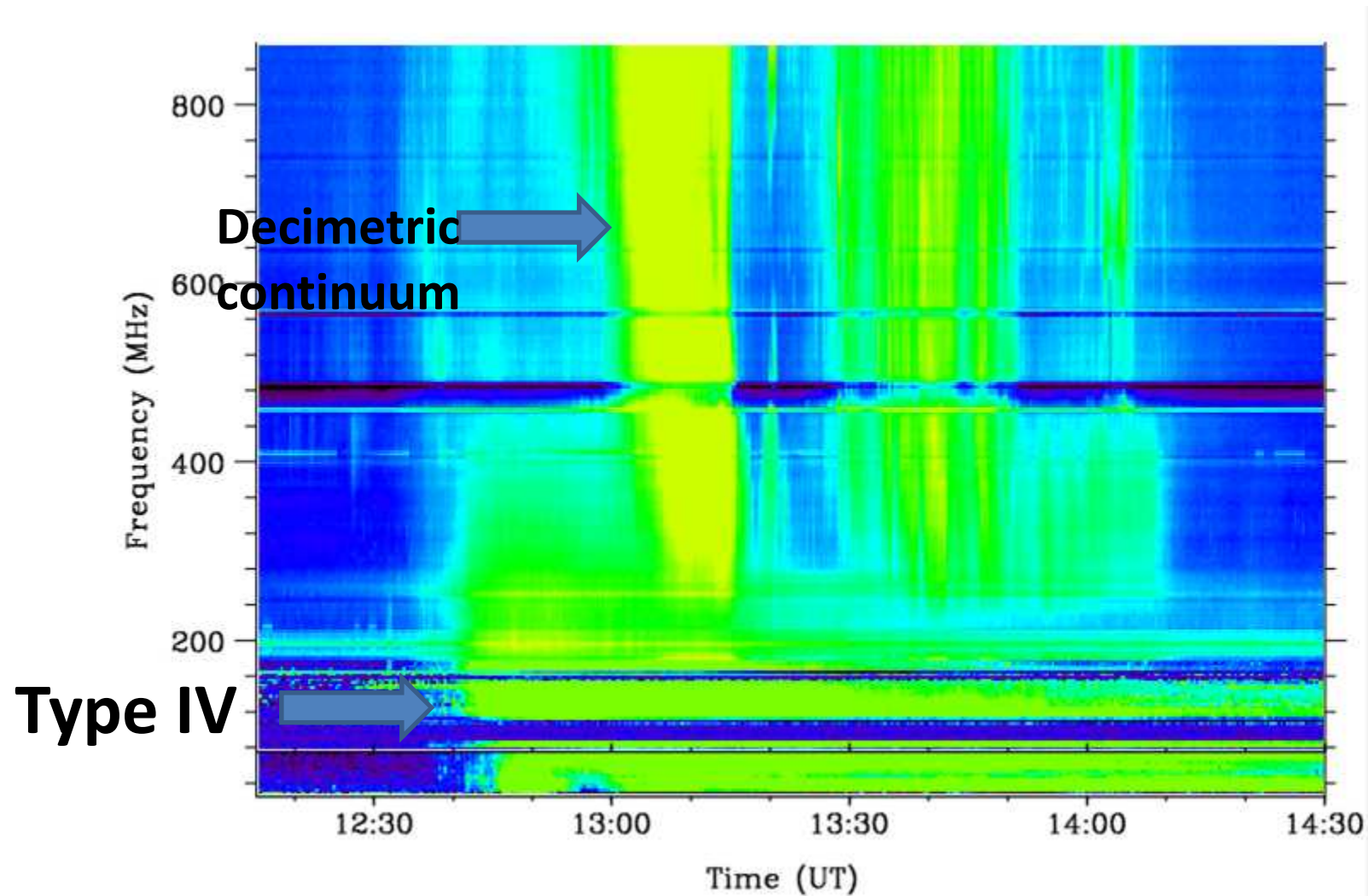
Moving Type IV(IVm) : Drift, Gyrosynchrotron, Plasma Radiation?



Stationary Type IV(IVs): No drift, Plasma Radiation?



Radio Dynamic Spectrum of 20110924



Important Physical Parameters

Brightness Temperature: T_b

Define a Temperature to indicate the intensity at a frequency that equals a **black body radiation**

Degree of Polarization: P

Stokes Parameters: I, Q, U, V , $Q^2 + U^2 + V^2 = I^2$ $I \sim$ Intensity

Circular Degree of Polarization $P = V/I$

Spectral Index: α

Given frequency f and radiative flux S , spectral index α satisfies

$$S = f^\alpha$$

Stationary Type IV

Common feature

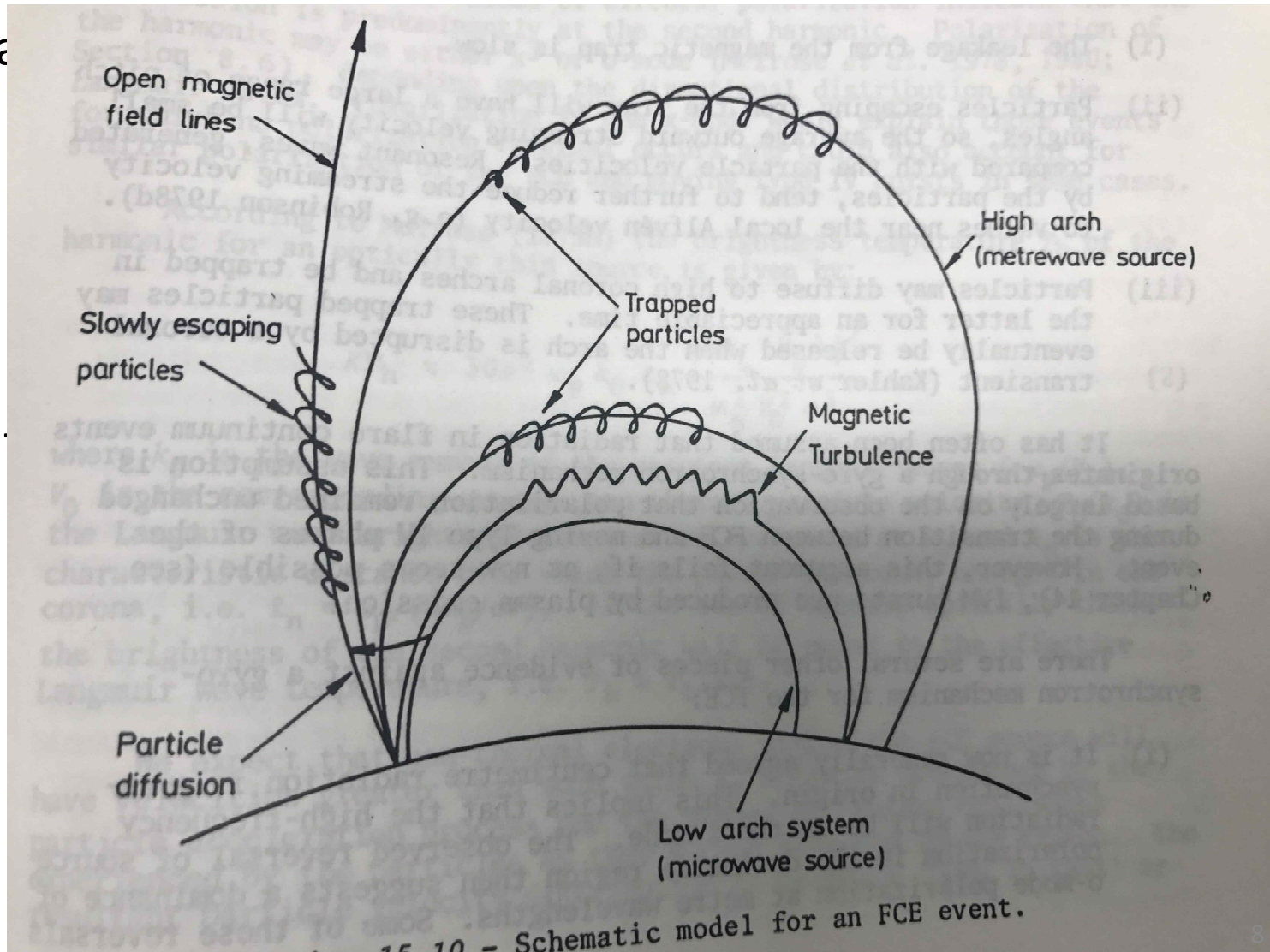
Polarization

Tb high

Continuous

Turn gradual

Mechanism



Radiation Mechanisms for Solar Radio Bursts

Synchrotron emission (fully relativistic)

Gyrosynchrotron radiation (partly relativistic)

Gyrosynchrotron: Cyclotron (non-relativistic) \times Synchrotron

Plasma Emission (fundamental, harmonic)

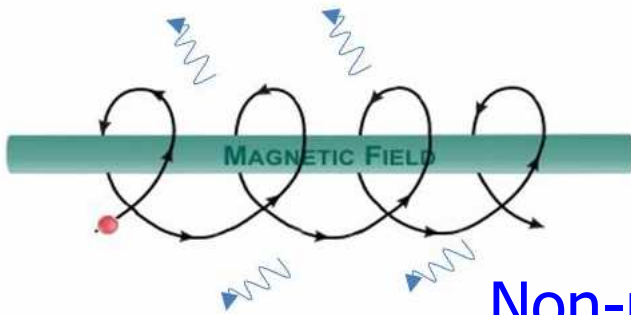
ECME (Electron Cyclotron Maser Effect)

Incoherent $\left\langle \text{-----} \right\rangle$ coherent

Tb low (max $\sim 10^9$)

Tb high

Gyrosynchrotron radiation



Non-relativistic: $\gamma-1 \ll 1$ (thermal)

cyclotron or gyroresonance radiation

Fully-relativistic: $\gamma-1 \gg 1$ (non-thermal)

synchrotron radiation

Mildly relativistic: $\gamma-1 \sim 1-5$ (thermal/non-thermal)

gyrosynchrotron radiation

Radiation frequency: gyrofrequency

Tb: less than 10^9

(Stewart, 1978)

Gyrosynchrotron radiation(incoherent)

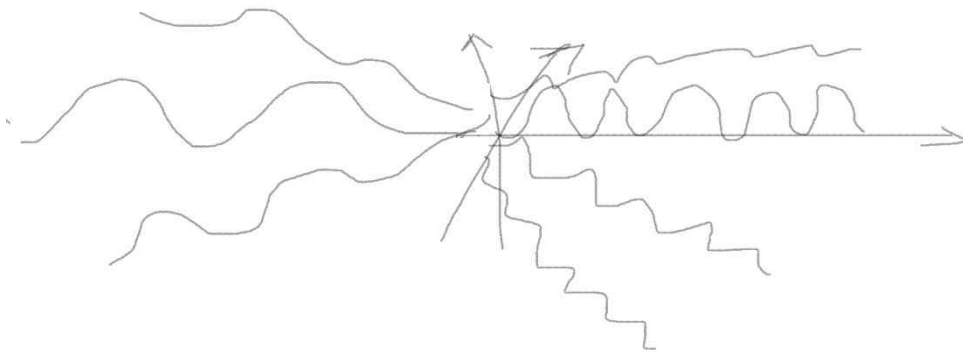
Cyclotron

Polarization:Any(elliptical)

Synchrotron emission(fully relativistic)

Polarization:

Linear Polarized (main reason:Doppler effect)



Gyrosynchrotron radiation:Any

Plasma Emission

$$f = 9000\sqrt{n_e} \quad f_{2h} \approx 2f_f$$

Fundamental : high p, Tb low usually

Second harmonic: low p, Tb high

Lakhina et al. 1985
(SoPh)

For most type I, II, IVs, V

Fundamental can be **AMPLIFIED** to $10^{16}K$

Second harmonic can be $10^{13}K$

Melrose
1975(SoPh)

Amplification of Plasma Emission

Only Fundamental :

In case of a *gap distribution*, fundamental can be AMPLIFIED to $10^{16}K$. Gap distribution requires:

(A) The distribution is isotropic.

(B) The number density of particles with $v \simeq v_0$ (e.g., in a range $\Delta v \simeq v_0$ about $v = v_0$) is sufficiently high for these particles to dominate over the thermal particles in the emission and absorption of Langmuir waves with $v_\phi \simeq v_0$.

(C) The number density of particles in the range $V_e \ll v \ll v_0$ is sufficiently low that absorption of Langmuir waves in the range $V_e \ll v_\phi \ll v_0$ is dominated by the thermal particles, while the emission of Langmuir waves in this range is dominated by the particles with $v \simeq v_0$.

Not possible for *power law* distributions

Melrose
1975(SoPh)

Electron Cyclotron Maser(ECM) Emission

Twiss (1958) first pointed out the possibility of the ECM instability to **directly amplify** the high-frequency electromagnetic waves at frequencies close to the **electron cyclotron frequency** and its **harmonics**

Electron Cyclotron Maser(ECM) Emission

Wu & Lee (1979) suggested electron distribution function, which is responsible for the ECM instability, is a so-called **loss-cone distribution** with $\frac{\partial f}{\partial v} > 0$ in a certain region around the loss cone and considering **relativistic effect**.

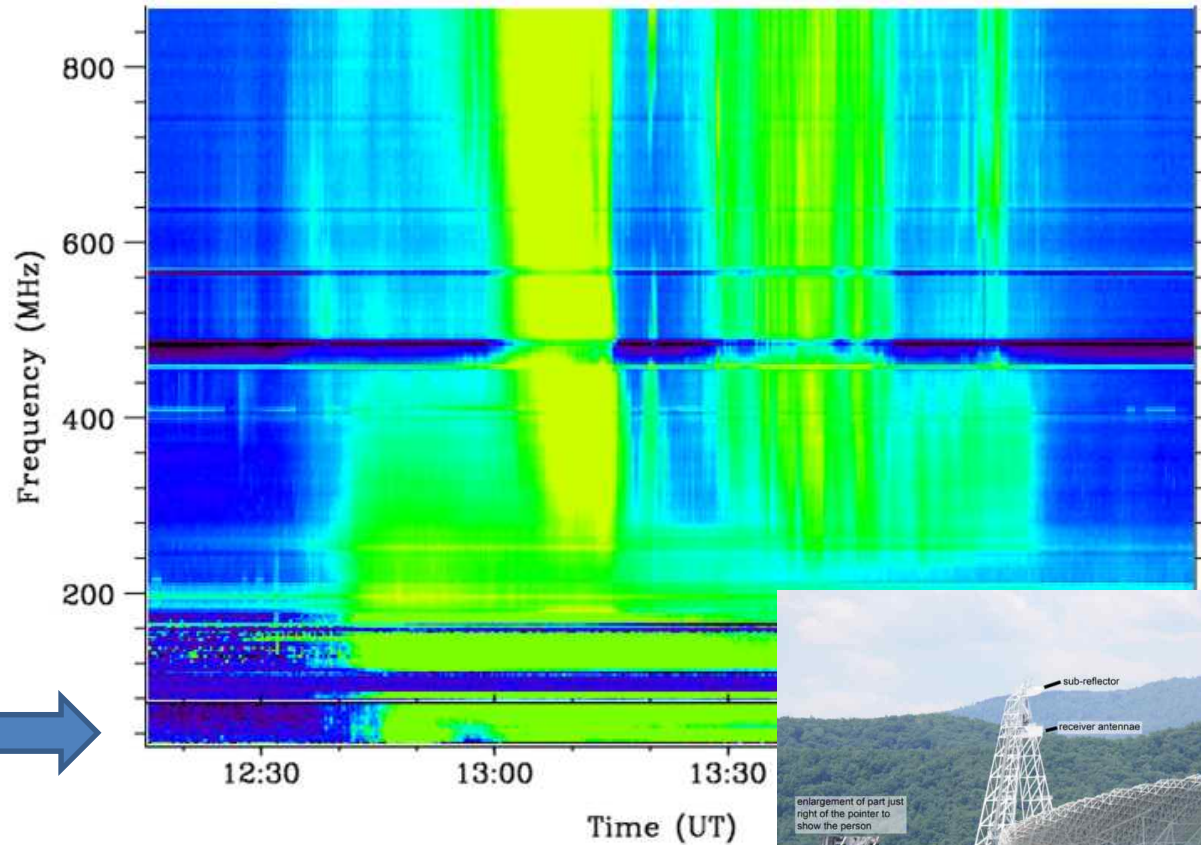
(f-distribution function)

This made it possible to explain radio bursts

Solar Radio Spectrometer

BLEN7M

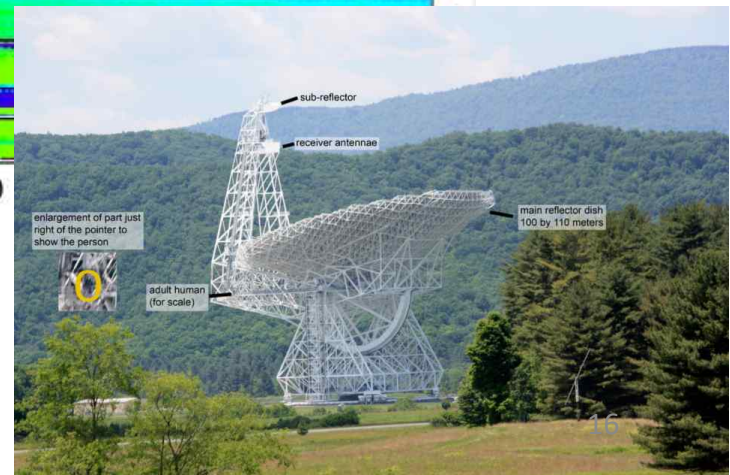
100-900 MHz



SGMR



25-75MHz



Solar Radioheliograph

Nançay RadioHeliograph(NRH,Fr) Tb, Stokes I, V

24 Sep 2011

Radio Monitoring

Movies

Starting hour : 12:00

Frame rate : 120 ▾ secs/frame

freq 150/164 MHz ▾

nb of images : 240

LOAD MOVIE

Home new day

Survey selection

Composite I : DAM
Composite IIa : DAM
Composite IIb : DAM
24h Waves/Swaves

Start Pause
Faster Slower
Reverse
Swing Mode: OFF
Rewind
Next Prev.
+10 -10



13:05:52 24SEP2011 150.9MHz

Frame: Displaying 33 of 96 Speed: 12 (frames/sec)



Others:

LOFAR(Netherland) MUSER(China) , GRAPH(India),UTR-2(Ukraine),MWA(Austalia)...

NoRH(17GHz&34GHz,Japan) SSRT(5.7GHz,Russia)...

Previous similar studies

My study will be conducted to

Get NRH&AIA combined data to locate the source

Tb,Polarization,particle distribution&Spectral index

Interpret radiation mechanism

previous:

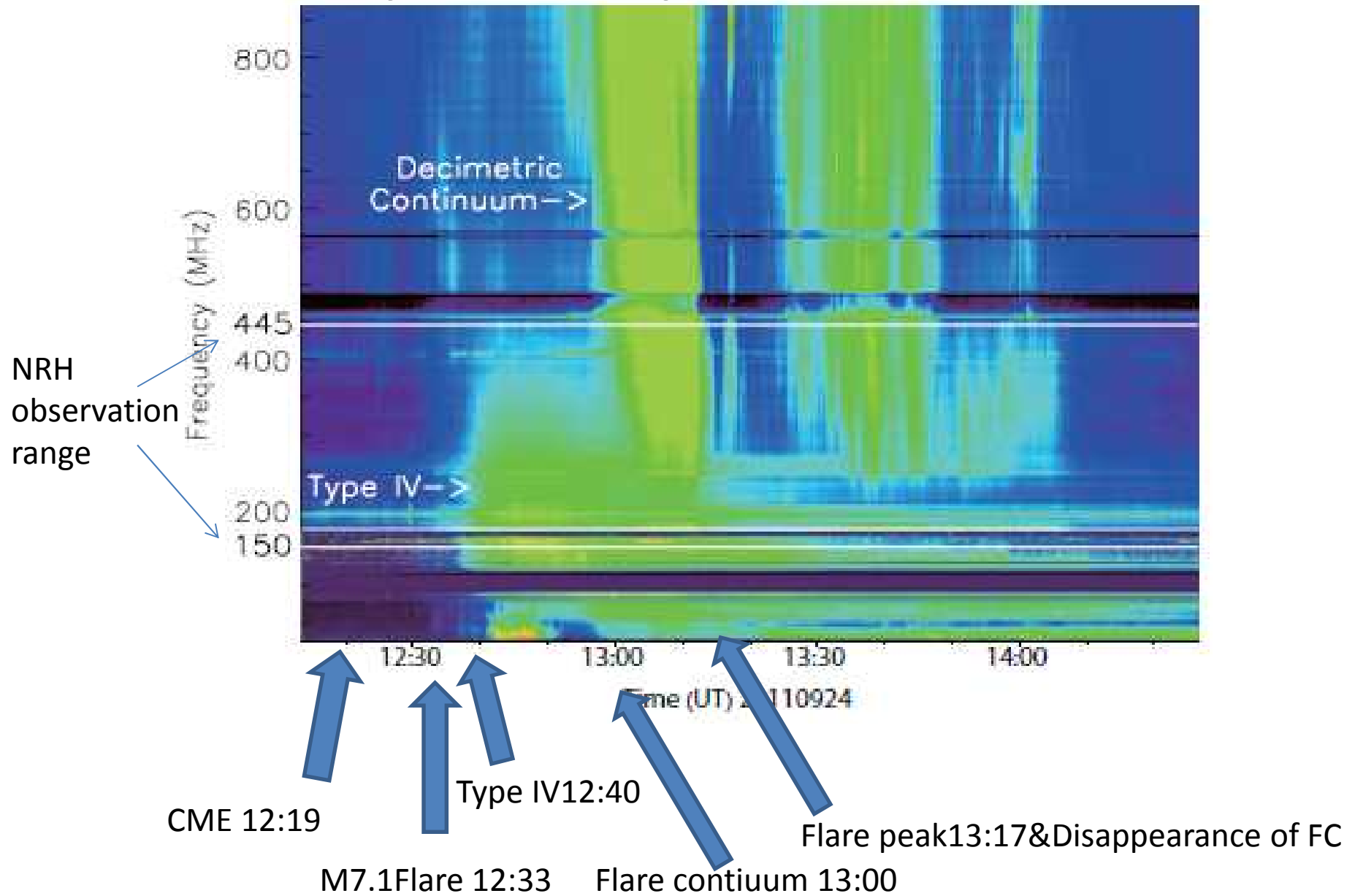
Bain, 2014 on a IVm, NRH&Stereo, Gyrosynchrotron

Vasanth, 2016 on a IVm, NRH&SDO, plasmoid,
eruptive hot channel

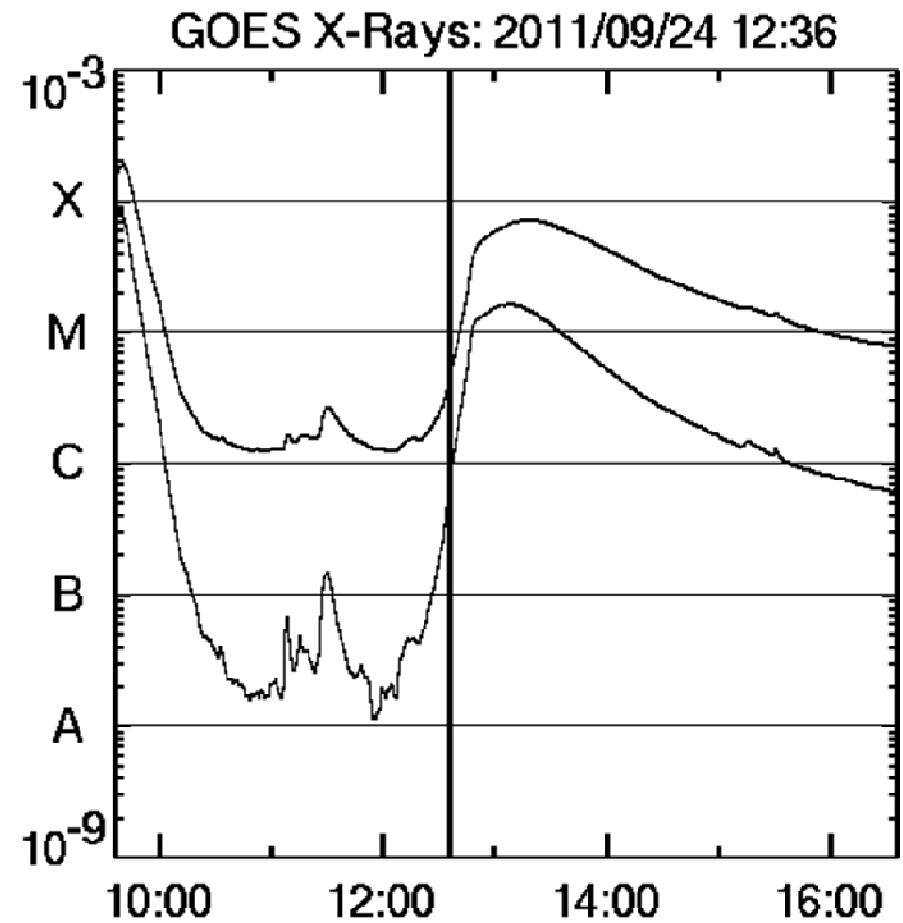
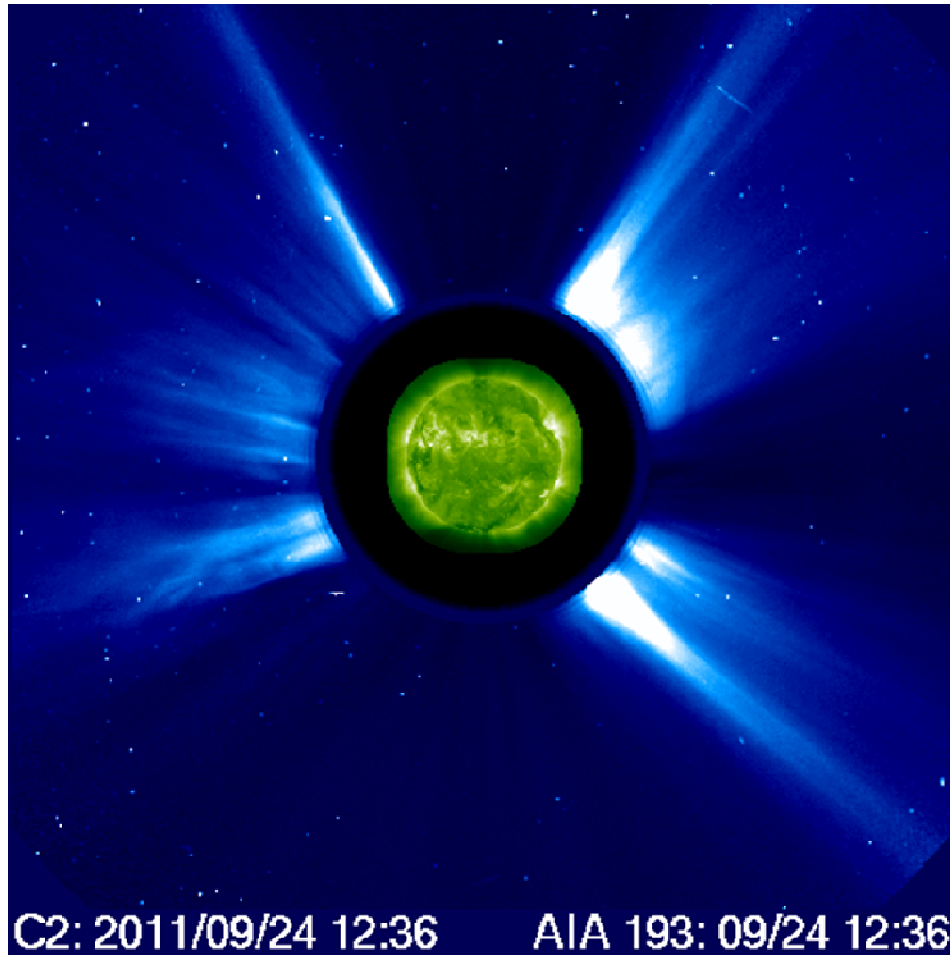
Koval, 2016 IVs, UTR-2&NRH&AIA&LASCO-C2, focus
on correlation with CME,no mechanism analysis

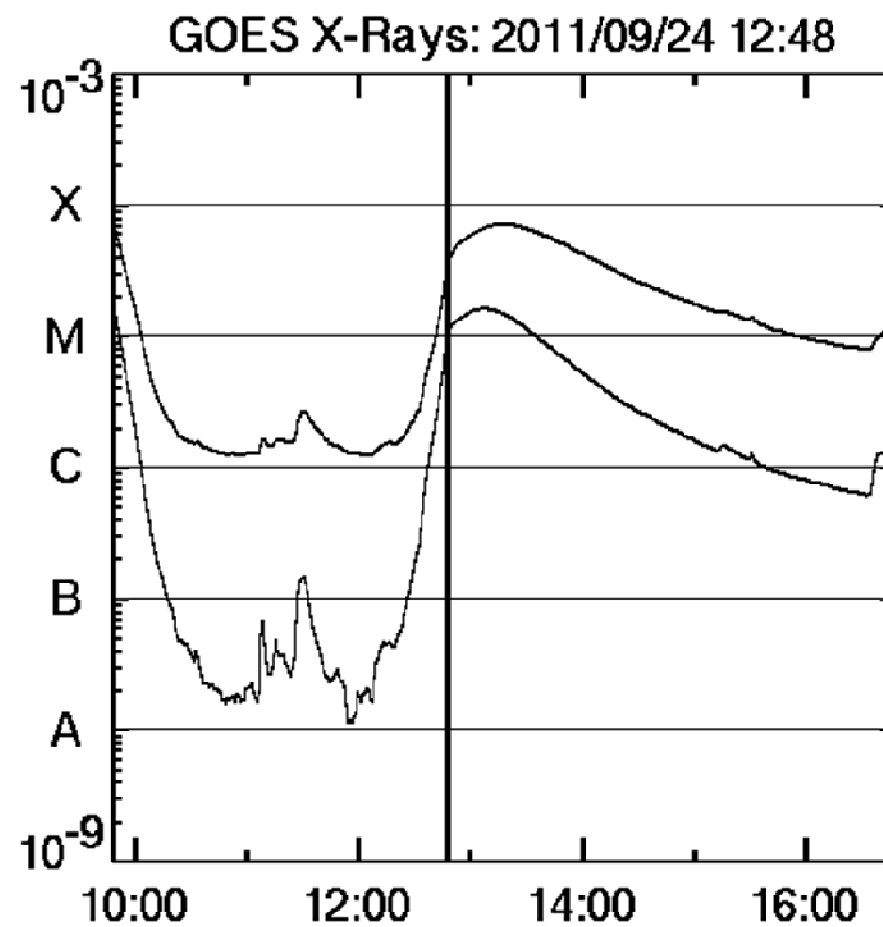
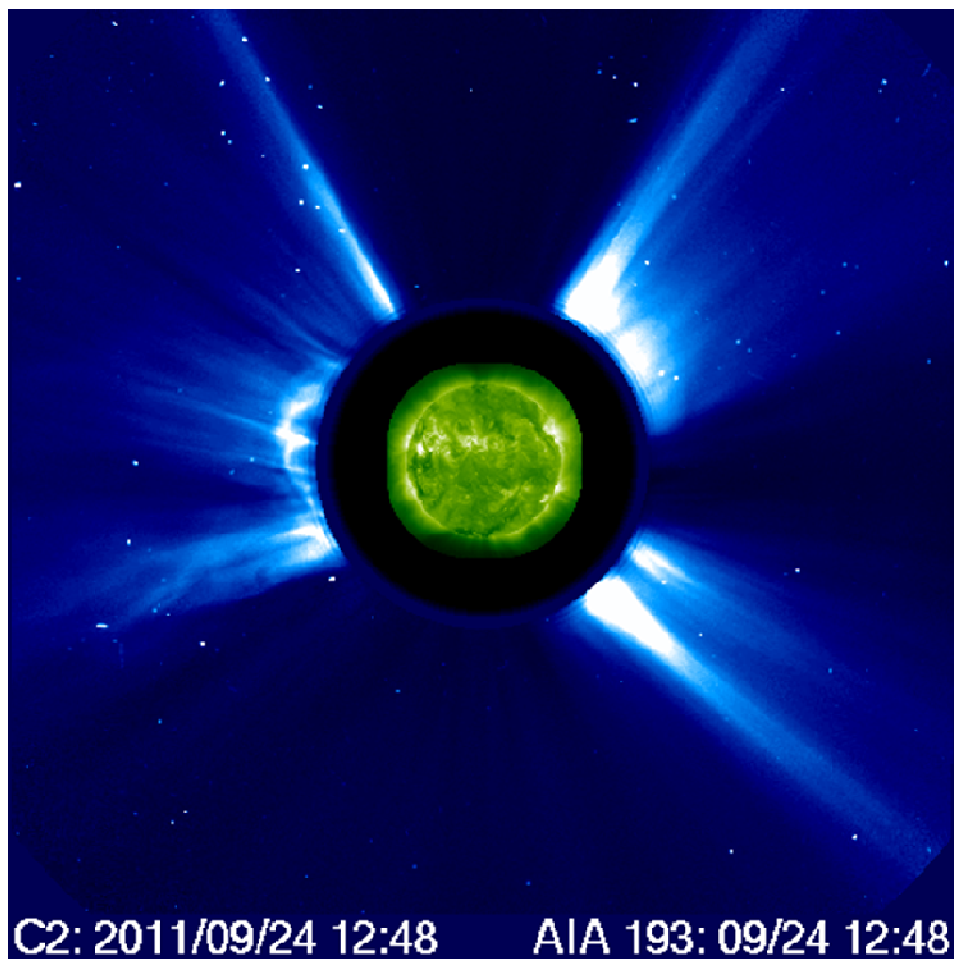
2 Study on Type IV

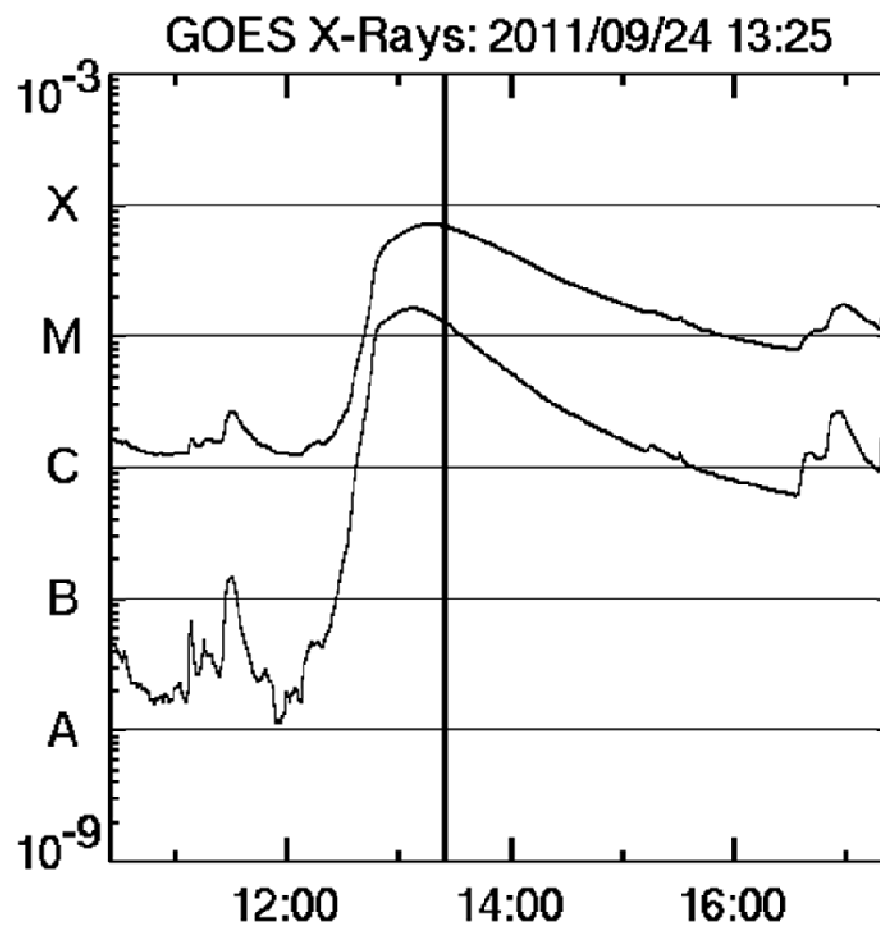
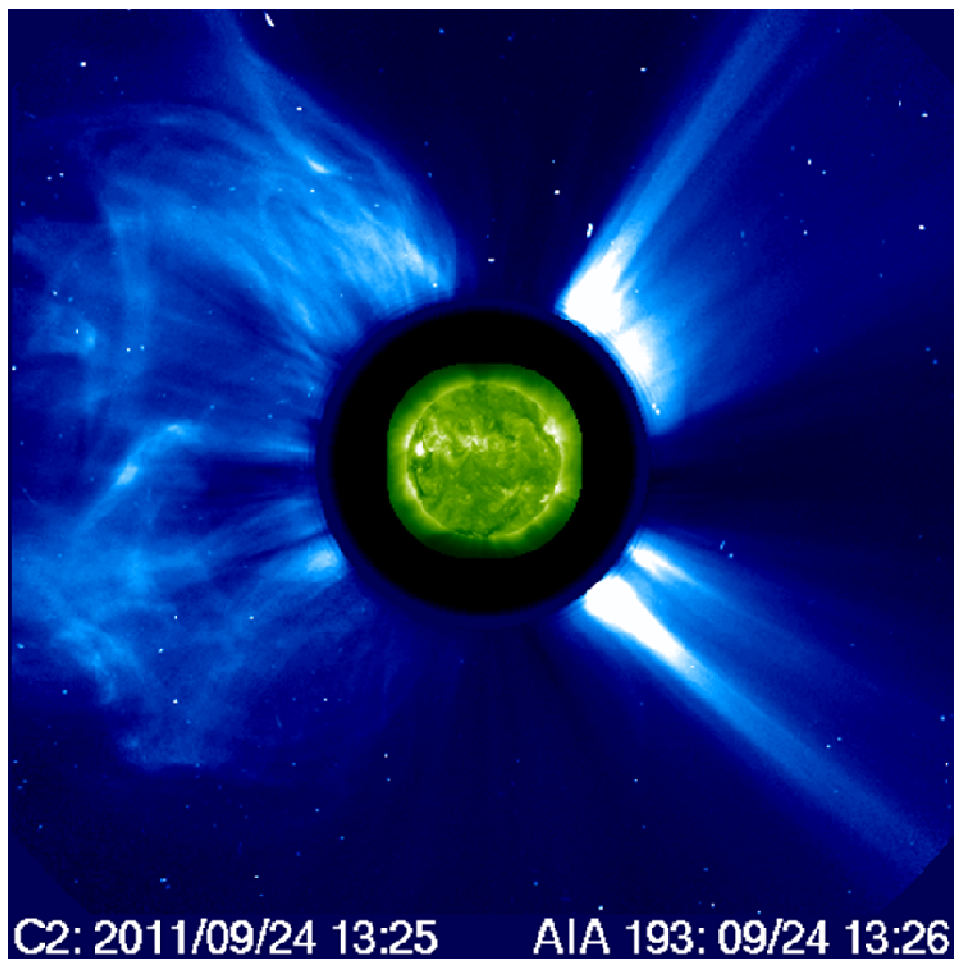
Radio Dynamic Spectrum of 20110924

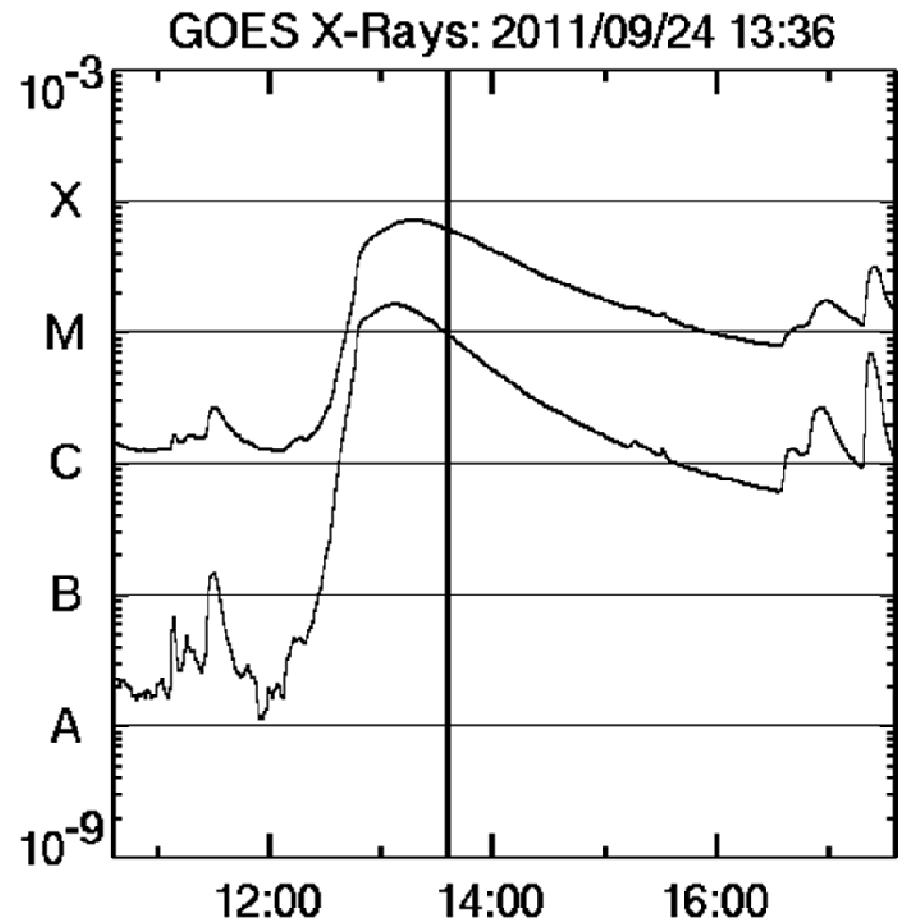
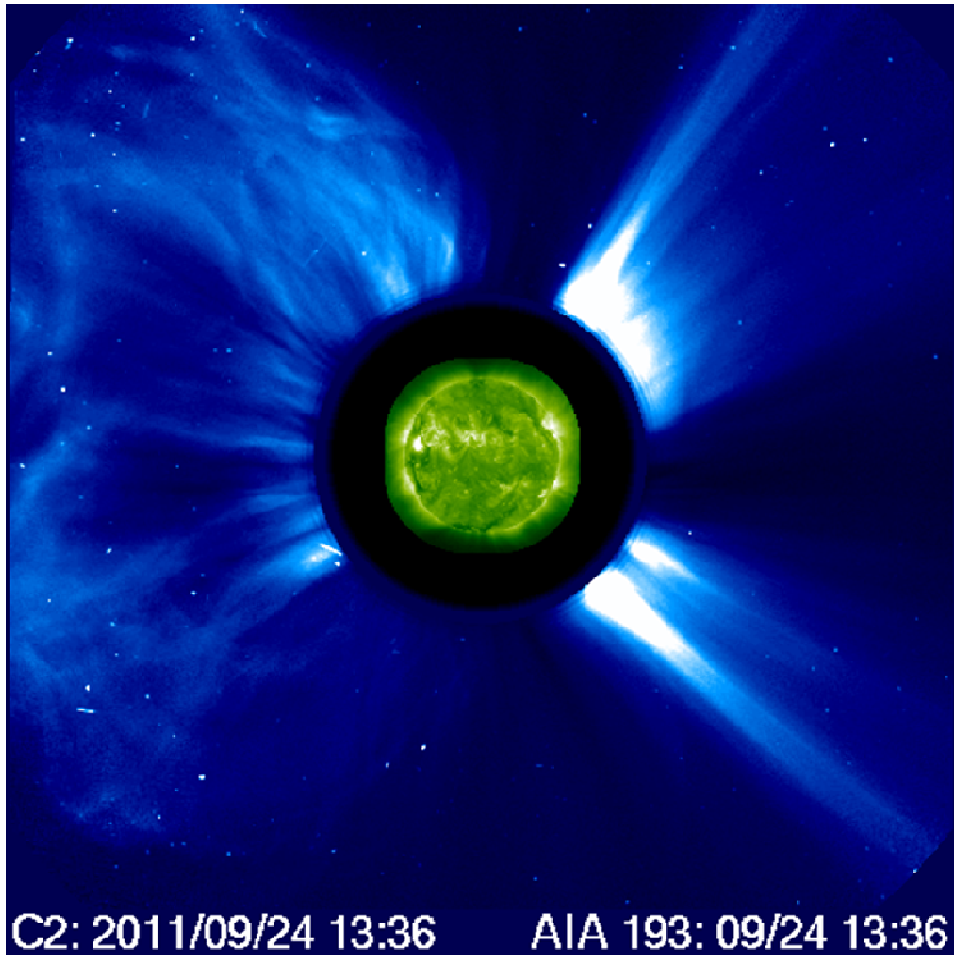


CME Before Type IV









AIA-NRH movie

94

1, stationary

131

2, sources on a open field
line, from the foot point,
Slowly escaping particles

171

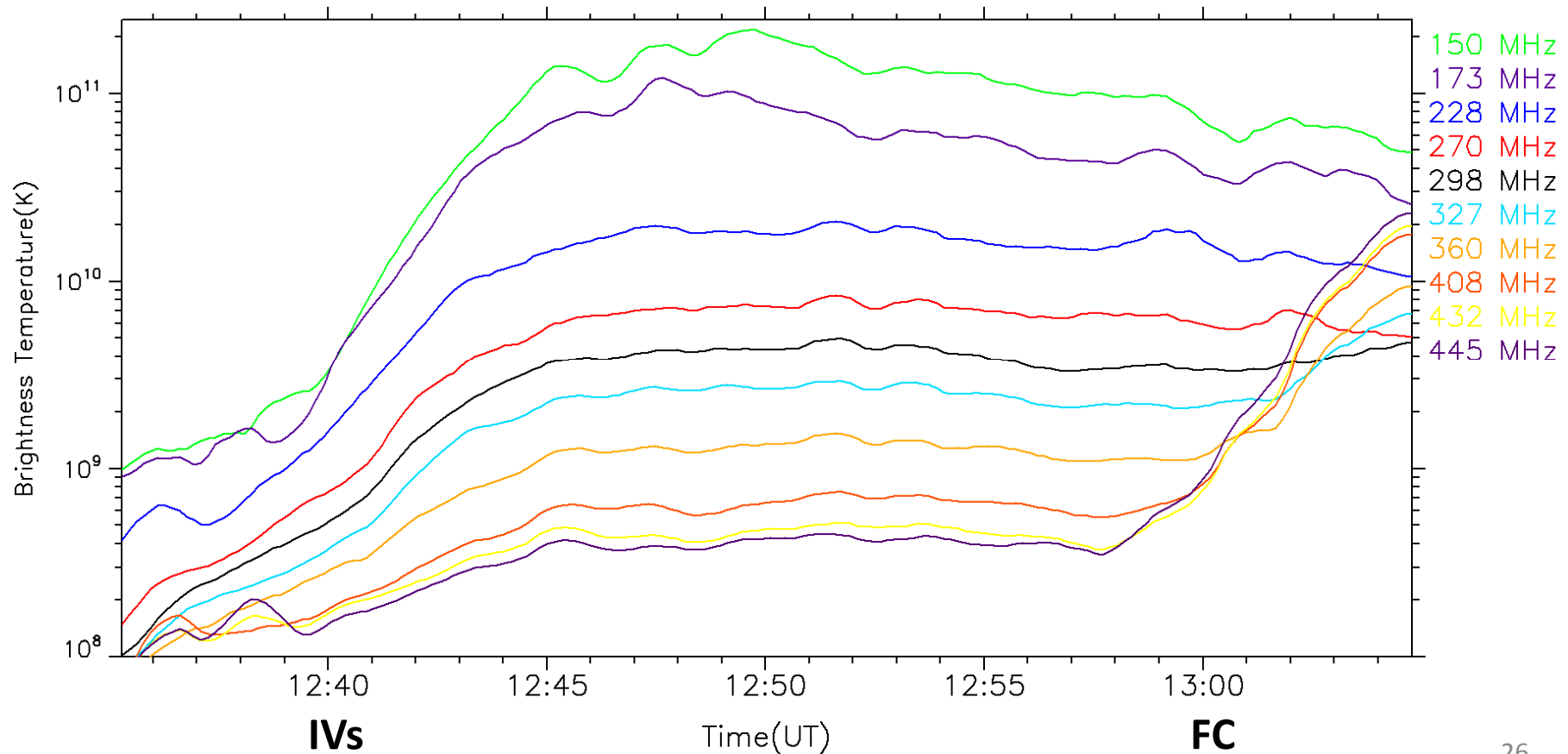
Brightness temperature calculation

Read directly from NRH fits file

Highest over 10^{11} K

frequency high, Tb low

Over 10^8 K

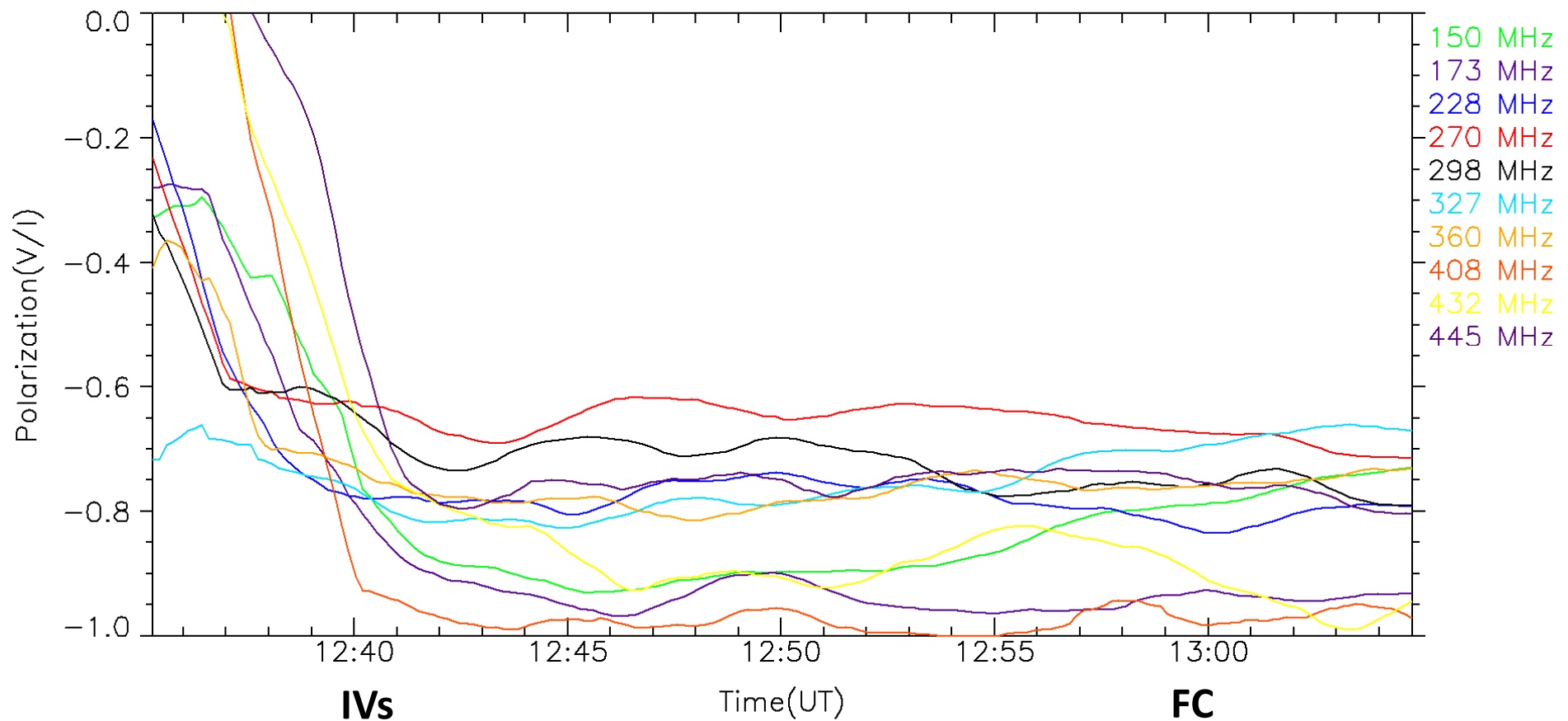


Degree of Polarization(circular) calculation

Read Stokes I and V from NRH fits file

$$P=V/I$$

Polarizations from 60% to 100% left-hand polarized



Flux-Frequency relation(distribution):

Read flux directly from NRH fits file

Type IV 12:40

Flare continuum 13:00

Always follows power-law distribution

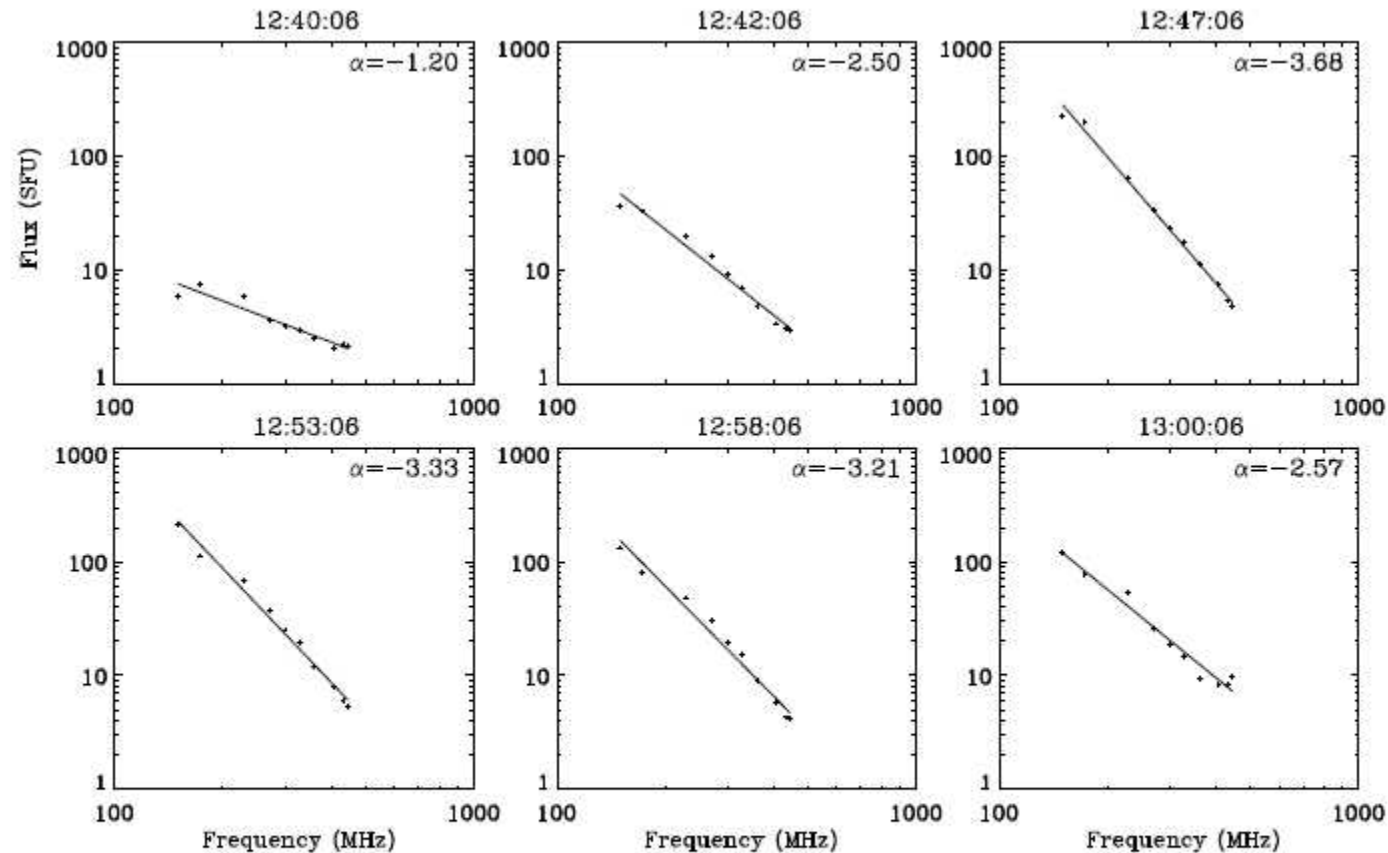


Figure 5. Linear fit of flux-frequency spectrum using NRH 10-frequency data from 12:40 UT to 13:00 UT. The dots are the flux data of 10 NRH frequencies in S.F.U. Spectral index (α) is shown on the upper right corner.

Spectral Index

Given frequency f and radiative flux S , spectral index α satisfies

$$S = f^\alpha$$

spectral index between -3 and -4

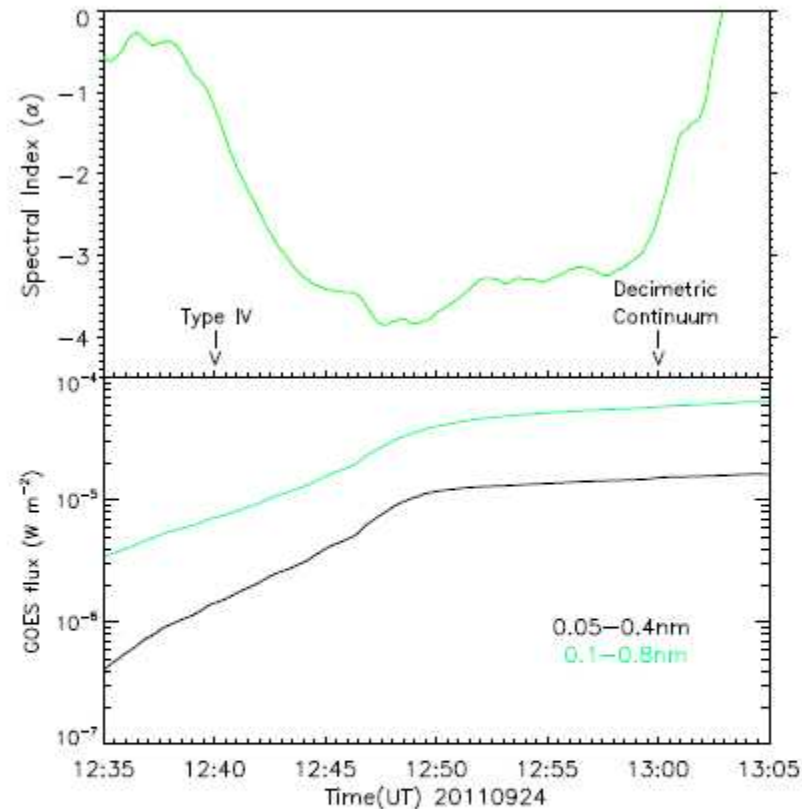


Figure 6. Temporal variation of the spectral index from 12:35 UT to 13:05 UT. GOES X-ray data are also plotted in the lower panel. The flare starts at 12:33 UT, and peaks at 13:17 UT, which means that the whole process happens during the rising phase of the flare. IVs occurs at 12:40 UT, and a decimetric continuum appears at 13:00 UT.

3 Summary and Radiation Mechanism Interpretation

Summary

* IVs with $T_b > 10^8 K$, sometimes $10^{11} K$;
high $P \sim 100\%$ left handed(-);
power law distribution,
with spectral index $\alpha -3 \sim -4$

* slowly escaping energetic particles along
magnetic field line, source from loop foot
point

Radiation Mechanism

Table 1. Characteristics of various radiation mechanisms and our observation. T_B stands for brightness temperature, q stands for degree of circular polarization, α means spectral index and δ is the particle energy distribution power law slope. We also list the observational parameters of IVs on Sep 24, 2011 for comparison.

Mechanism	T_B	Polarization ($ q $)	Spectral Index	Reference
Cyclotron	$<10^9$ K	Any	$\alpha \neq f(\delta)$	Dulk (1985)
Synchrotron	$<10^9$ K	0% (Linear)	power law, $\alpha \propto \delta$	Dulk (1985)
Gyrosynchrotron	$<10^9$ K	Any	$\alpha \neq f(\delta)$	Robinson (1978)
Fundamental Plasma	$\sim 10^9$ K / $< 10^{16}$ K	$\sim 100\%$ / $\sim 0\%$	$\alpha \neq f(\delta)$ or $\alpha \propto \delta^1$	Melrose (1975)
2nd Harmonic Plasma	$< 10^{13}$ K	$< 10\%$		Melrose (1975)
ECM Emission	$\geq 10^{10}$ K	$\sim 100\%$	power law	Winglee (1985)
Our Observation				
IVs of 20110924	10^{11} K	60-100%	power law	

¹For relativistic electrons

T_b is high, so there should be considerable relativistic effect. For relativistic electrons, for plasma radiation, $\alpha \propto \delta$ (Kaplan and Tsytovich, 1969). fundamental plasma radiation can't be amplified by power law electrons

ECM Emission!

Radiation Mechanism

ECM emission usually exhibits **small time-scale spikes**, which is **less than 0.1s** (Aschwanden, 2004).

However, we reiterate that **long duration continuum** has **also** been **explained by ECM emission**(Winglee and Dulk, 1986).

Moreover, Zhao et al. (2013) has developed a theory to **explain** the **type I radio burst** with the **ECM** mechanism, which is also **long duration continuum**.

Type I Burst(Noise Storm)

Common features:

Long duration

Continuum emission

Polarization high~ 100%

Tb High

Conclusion

IVs

Source region along magnetic field line(open or high lying loop)

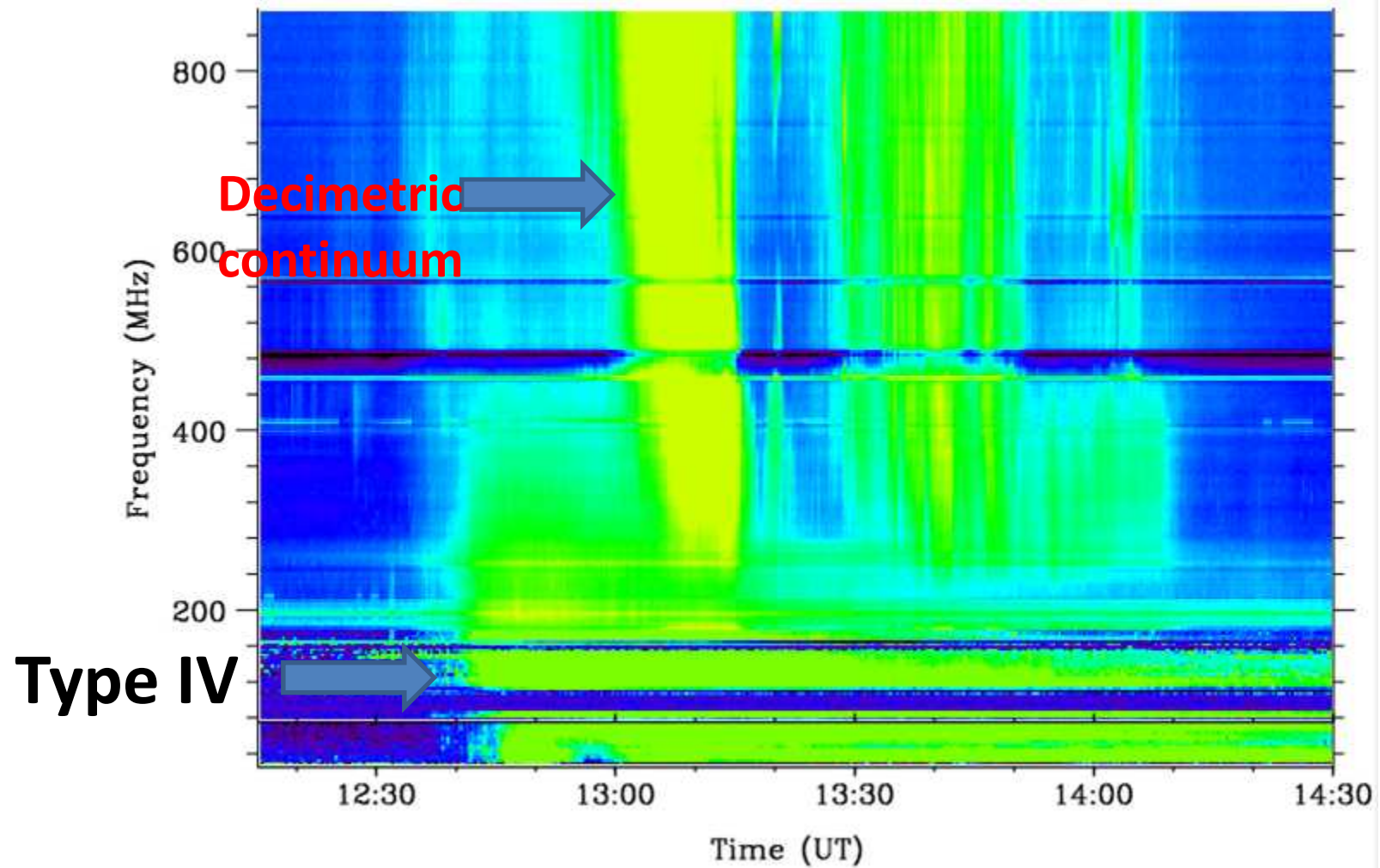
Tbis high $10^8 \sim 10^{11} \text{K}$

Degree of circular polarization $-60 \sim -100\%$

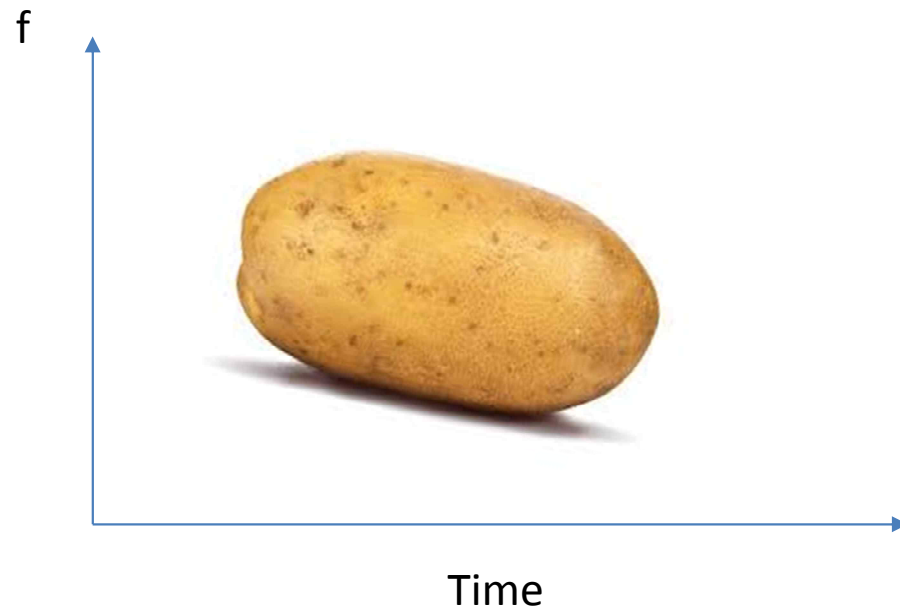
Power law, spectral index $-3 \sim -4$

ECM emission mechanism

Discussion



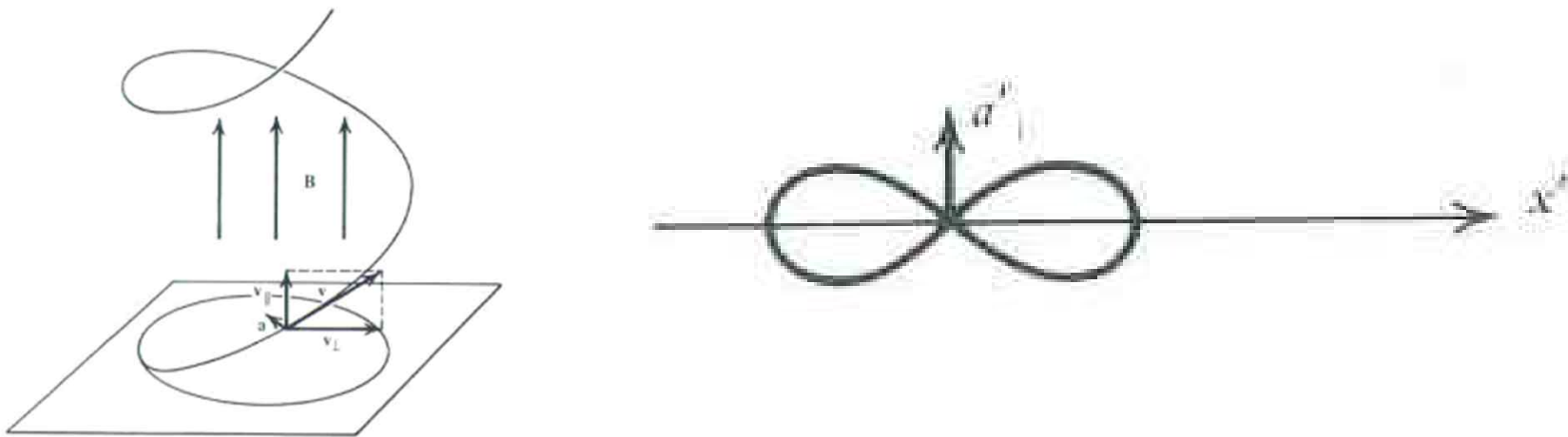
Sorry for the long post



Gyroemission

Gyroemission is due to the acceleration experienced by an electron as it gyrates in a magnetic field due to the Lorentz force. The acceleration is perpendicular to the instantaneous velocity of the electron.

When the electron velocity is nonrelativistic ($v \ll c$ or $\gamma - 1 \ll 1$) the radiation pattern is just the dipole pattern.



Non-relativistic: $\gamma^{-1} \ll 1$ (thermal)

cyclotron or gyroresonance radiation

Mildly relativistic: $\gamma^{-1} \sim 1-5$ (thermal/non-thermal)

gyrosynchrotron radiation

Ultra-relativistic: $\gamma^{-1} \gg 1$ (non-thermal)

synchrotron radiation

Radiation frequency: gyrofrequency

Tb: less than 10^9
(Stewart, 1978)

Gyrosynchrotron radiation(incoherent)

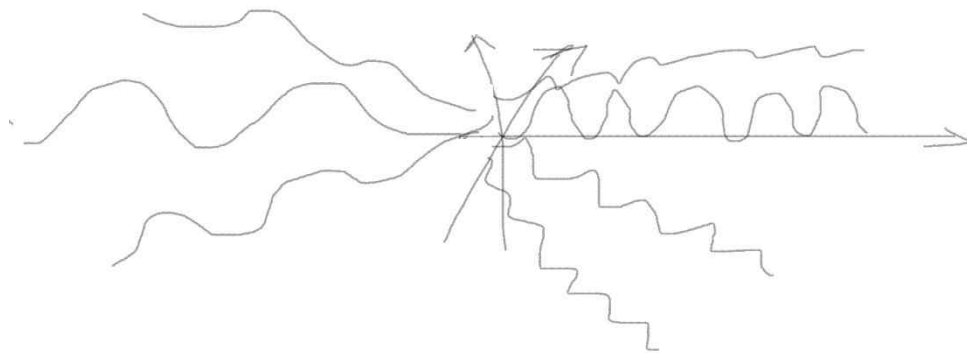
Cyclotron non-relativistic

Polarization:Any(elliptical)

Synchrotron emission(fully relativistic)

Polarization:

Linear Polarized (main reason:Doppler effect)



Gyrosynchrotron radiation:Any

Plasma radiation

Plasma oscillations (**Langmuir waves**) are a **natural mode** of a plasma and can be excited by a variety of mechanisms.

In the Sun's corona, the propagation of **electron beams** and/or **shocks** can excite plasma waves.

These are converted from **longitudinal** oscillations to **transverse** oscillation through nonlinear wave-wave interactions.

The resulting transverse waves have frequencies near the **fundamental** or **harmonic** of the local electron plasma frequency: i.e., ν_{pe} or $2\nu_{pe}$.

Plasma radiation

Plasma radiation is therefore thought to involve several steps:

Fundamental plasma radiation

- A process must occur that is unstable to the production of Langmuir waves
- These must then scatter off of thermal ions or, more likely, low-frequency waves (e.g., ion-acoustic waves)

$$\omega_L + \omega_S = \omega_T$$

and

$$k_L + k_S = k_T$$

coalescence

or

$$\omega_L = \omega_S + \omega_T$$

$$k_L = k_S + k_T$$

decay

Plasma radiation

Plasma radiation is therefore thought to involve several steps:

Harmonic plasma radiation

- A process must occur that is unstable to the production of Langmuir waves
- A **secondary spectrum** of Langmuir waves must be generated
- Two Langmuir waves can then coalesce

$$\omega_L^1 + \omega_L^2 = \omega_T$$

and

$$k_L^1 + k_L^2 = k_T \ll k_L$$

$$\omega_T \approx 2\omega_L$$

$$k_L^1 \approx -k_L^2$$