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(1)

# Radiative Transfer Equation - Recap

$$\mu \frac{dI}{d\tau} = I - S$$

To be complete

$$\mu \frac{dI_{\nu}(\tau, \mu)}{d\tau_{\nu}} = I_{\nu}(\tau, \mu) - S_{\nu}(\tau, \mu) \quad \dots (4.10)$$

$$d\tau_{\nu} = -\kappa_{\nu} dz \quad ; \quad dz = \mu ds = \cos\theta ds$$

$$S_{\nu} = \frac{j_{\nu}}{\kappa_{\nu}} = \text{source function}$$

General solution

$$I_{\nu}(\tau, \mu) = e^{-(\tau_0 - \tau)/\mu} I_{\nu}(\tau_0, \mu) + \int_{\tau}^{\tau_0} e^{-(\tau - t)/\mu} \frac{S_{\nu}(t)}{\mu} dt \quad \dots (4.11)$$

# Stellar interior

$$\boxed{S(\tau) = B(\tau) + (\tau - \tau_0) \left( \frac{\partial B}{\partial \tau} \right)_{\tau}} \quad \dots (4.14)$$

$$I(\tau, \mu > 0) = B(\tau) + \mu \left( \frac{\partial B}{\partial \tau} \right)_{\tau} \quad \dots (4.15)$$

$$I(\tau, \mu < 0) \approx B(\tau) + \mu \left( \frac{\partial B}{\partial \tau} \right)_{\tau} \quad \dots (4.16)$$

$$F(\tau) = 2\pi \int_{-1}^1 I(\mu) \mu d\mu$$

$$\boxed{F(\tau) = \frac{4\pi}{3} \frac{\partial B(\tau)}{\partial \tau}} \quad \dots (4.17)$$

$$L(r) = 4\pi r^2 F(r) \quad \dots (4.18)$$

## # Measure of Radiation Anisotropy

$$\text{Anisotropy} = \frac{F}{I} = \frac{\text{net flux}}{\text{Intensity}} = \frac{\frac{4\pi}{3} \frac{\partial B}{\partial t}}{B_{\text{CMB}}}$$

Example: the Sun

$$L = 4\pi R^2 F$$

$$L_0 = 4 \times 10^{33} \text{ erg s}^{-1}$$

$$R_0 = 6.96 \times 10^{10} \text{ cm}$$

$$F = \frac{L}{4\pi R^2} = \frac{4 \times 10^{33}}{4\pi \times (6.96 \times 10^{10})^2} = 2 \times 10^{10} \text{ erg cm}^{-2} \text{ s}^{-1}$$

$$B_{\text{CMB}} = \frac{\sigma T_0^4}{\pi}$$

$$T_0 \sim 10^4 \text{ K}$$

$$\sigma = 5.67 \times 10^{-5} \text{ erg cm}^{-2} \text{ K}^{-4} \text{ s}^{-1}$$

$$= 2 \times 10^{23} \text{ erg cm}^{-2} \text{ s}^{-1} \text{ sr}^{-1}$$

$$\text{Anisotropy} = \frac{2 \times 10^{10}}{2 \times 10^{23}} = 10^{-13}$$

in stellar interior,

Therefore, the luminosity is caused by a "small" anisotropy of the radiation field, which is further caused by the temperature gradient.

# The Diffusion Equation (CH 4.2)

~ radiation flux in classical diffusion format

$$F = \frac{4\pi}{3} \frac{\partial B(r)}{\partial r}$$

$$\Rightarrow F = -D_{\text{rad}} \frac{dT}{dr}$$

Fick's Law  
of Diffusion

# Diffusion:  $D_{\text{rad}}$ : Diffusion coefficient of radiation  
from area of high density to area of low density

$$\frac{\partial B(r)}{\partial r} = \frac{\partial B(r)}{\partial T} \frac{\partial T}{\partial r}$$

$$B(r) = \frac{cT^4}{\pi} = \frac{c}{4\pi} aT^4 ; \quad \frac{\partial B(r)}{\partial T} = \frac{c}{\pi} aT^3$$

$$\partial r = -\kappa \rho \partial r$$

$$\Rightarrow F(r) = -\frac{4ac}{3} \frac{1}{\kappa \rho} T^3 \frac{dT}{dr} \quad \text{--- (4-23)}$$

Therefore  $D_{\text{rad}} = \frac{4ac}{3} \frac{1}{\kappa \rho} T^3$

Another form:  $F(r) = -\frac{c}{3\kappa \rho} \frac{d(c a T^4)}{dr}$

$U = aT^4$ : radiation energy density

$\frac{1}{\kappa \rho} = \lambda$ : photon mean free path

$$F(r) = -\frac{1}{3} c \lambda \frac{dU}{dr}$$

$U = C_v T$ : For gas

$$F(r) = -\frac{1}{3} C_v c \lambda \frac{dT}{dr}$$

which is similar to

classical thermal conduction formula

$$F = -\frac{1}{3} C_v v_e \lambda \frac{dT}{dr} ; \quad v_e: \text{mean electron velocity}$$

$\lambda$ : mean electron free path